

## Modified Approach to Calculate Water Influx (Unsteady State)

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### طريقة معدلة لحساب تدفق الماء غير المستقر

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تم في هذه الورقة وباستعمال طرق مختلفة<sup>[4,3,2]</sup> حل مثالين لحساب تدفق الماء غير المستقر من كتاب أساسيات الهندسة النفطية<sup>[1]</sup> بهدف توضيح أن معادلات Marsal مازالت في حاجة إلى بعض التحوير في الثوابت حتى تتوافق مع الحل الأصلي لـ Van Everdingen و Hurst. إن هذه التعديلات غير ثابتة ولكنها تعتمد على المدى لقيم  $t_D$ ،  $t_{ED}$ .

**Abstract:** In this paper two examples for the calculation of water influx from Ref<sup>[1]</sup> are resolved using different techniques (Van Everdingen and Hurst<sup>[2]</sup>, Fetkovich<sup>[3]</sup>, Marsal<sup>[4]</sup> in order to demonstrate that Marsal Equations need some modification on the constants in order to be more consistent with the original solution of Van Everdingen and Hurst<sup>[2]</sup>. These modifications are not constant but are function of the range of  $r_{eD}$  and  $t_D$ .

### INTRODUCTION

A great number of papers were written about the subject of water influx calculations in unsteady state for radial and linear flow; so there is no need to go in detail about the theory of this subject .

The unsteady state influx theory of Van Everdingen and Hurst is used to calculate the cumulative water influx, under practically all circumstances, for radial and linear aquifers. Unfortunately, it has the disadvantage that the calculations are rather tedious due to the complexity of superposing solutions for each time step, also it requires tables of  $We_D$  and  $t_D$ .

Fetkovich tried to avoid or remove the necessity of superposition in the Van Everdingen and Hurst method. However it is used for limited aquifers only. Also it is initially still necessary to apply the unsteady state influx theory of Van Everdingen and Hurst for the first few time steps.

Marsal has shown that the Van Everdingen and Hurst tables of  $We_D$  and  $t_D$  can be fairly well approximated by the infinite radial unit function up to a dimensionless time of  $t_D = 0.4 (r_{eD} - 1)^2$ . He used analytical expression to calculate the dimensionless water influx ( $We_D$ ).

The modified approach used almost the same analytical expression of Marsal but with slight modifications on the constants and there are some constrains on the values of  $r_{eD}$  and  $t_D$ .

### THE MODIFIED APPROACH

Marsal has shown that the Van Everdingen and Hurst tables of  $We_D$  and  $t_D$  can be fairly well approximated by the infinite radial unit function up to a dimensionless time of  $t_D = 0.4 (r_{eD} - 1)^2$  .....(1)

The approximating function can be written in the following form:

$$We_D = 0.5 (r_{eD}^2 - 1) [1 - \exp (- 2 t_D / j^* )].....(2)$$

Where :

$$j^* = r_{eD}^4 \ln (r_{eD}) / (r_{eD}^2 - 1) + 0.25 (1 - 3 r_{eD}^2) .....(3)$$

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$$t_D = 2.309kt / \Phi \mu_w c_r r_o^2 \dots\dots\dots(4)$$

$$r_{eD} = r_e / r_o \dots\dots\dots(5)$$

The dimensionless time and aquifer constant for finite radial aquifers are defined as for infinite radial aquifers.

Applications of the above equations on the examples in Ref [1] gave some deviation from the results obtained using method of Van Ever and Hurst, especially for large value of  $r_{eD}$ . So trial and error was used on the constants in the above equations and some constrains were put on the values of  $r_{eD}$  and  $t_D$  as follow:

For  $r_{eD} \leq 6$  then

$$j^* = r_{eD}^4 \ln(r_{eD}) / (r_{eD}^2 - 1) + 0.245 (1 - 3 r_{eD}^2) \dots\dots(6)$$

For  $r_{eD} > 6$  then

$$j^* = r_{eD}^4 \ln(r_{eD}) / (r_{eD}^2 - 1) + 0.3071 (1 - 3 r_{eD}^2) \dots\dots(7)$$

For  $t_D \leq 10$  then

$$W_{eD} = 0.5 (r_{eD}^2 - 1) [ 1 - \exp(- 2.2 t_D / j^*) ] \dots\dots(8)$$

For  $10 < t_D \leq 20$  then

$$W_{eD} = 0.5 (r_{eD}^2 - 1) [ 1 - \exp(- 2.15 t_D / j^*) ] \dots\dots(9)$$

For  $t_D > 20$  then

Table 1. Input data

Time (years)	$t_D$	Pressure at o.w.c (psia)	$W_{eD}(r_{eD} = 10)$	$W_{eD}(r_{eD} = 5)$
0	0.00	2720 ( $p_i$ )	0.00	0.00
1	5.67	2500	4.95	4.88
2	11.34	2290	8.12	7.46
3	17.01	2109	10.90	9.10
4	22.68	1949	13.50	10.09
5	28.35	1818	15.90	10.83
6	34.02	1702	18.10	11.27
7	39.69	1608	20.20	11.52
8	45.36	1535	22.20	11.69
9	51.03	1480	24.00	11.81
10	56.70	1440	25.70	11.89

$h = 100$  ft,  $k = 200$  md,  $r_o = 9200$  ft.,  $c_w = 3 \times 10^{-6}$  psi<sup>-1</sup>,  $c_r = 4 \times 10^{-6}$  psi<sup>-1</sup>,  $f = 140^\circ / 360^\circ$ ,  $\Phi = 0.25$ ,  $\mu_w = 0.55$  cp.

$$W_{eD} = 0.5 (r_{eD}^2 - 1) [ 1 - \exp(- 1.875 t_D / j^*) ] \dots\dots(10)$$

Examples

Two examples were taken from reference [1]. The input data are shown in Table 1. The four methods were applied using a short computer program (available on request) to calculate the cumulative water influx .The results are shown in Tables 2 and 3.

Table 2. Results of cumulative water influx calculated by different methods x10<sup>6</sup>bbl  $r_{eD} = 5$

Time Year	$t_D$	Van and Hurst	Fekvich	Marsal	Modified
1	5.67	3.77	3.93	3.56	3.78
2	11.34	12.84	13.55	12.43	13.05
3	17.01	24.09	25.51	23.71	24.51
4	22.68	35.67	37.98	35.72	36.46
5	28.35	47.25	49.95	47.43	47.70
6	34.02	58.02	60.88	58.33	58.25
7	39.69	67.76	70.64	68.15	67.84
8	45.36	76.23	79.00	76.68	76.24
9	51.03	83.36	85.93	83.84	83.32
10	56.70	89.21	91.47	89.77	89.11

Table 3. Results of cumulative water influx calculated by different methods x10<sup>6</sup>bbl  $r_{eD} = 10$

Tim Year	$t_D$	Van and Hurst	Fekvich	Marsal	Modified
1	5.67	3.82	2.72	2.65	3.25
2	11.34	13.46	10.35	10.09	12.20
3	17.01	26.44	21.88	21.35	25.46
4	22.68	41.92	36.46	35.59	40.94
5	28.35	59.10	53.31	52.07	57.61
6	34.02	77.46	71.75	70.13	76.00
7	39.69	96.66	91.27	89.92	95.40
8	45.36	116.14	111.30	108.92	115.37
9	51.03	135.46	131.35	128.62	135.37
10	56.70	154.29	151.06	148.02	154.98

CONCLUSION

The results obtained using the suggested modifications of the Marsal Equations indicate clearly that the the Marsal constants need to be varied based on the values of  $r_{eD}$  and  $t_D$  .It is suggested that the Marsal Equations be re –written in the following form:

$$J^* = r_{eD}^4 \ln(r_{eD}) / (r_{eD}^2 - 1) + C_1 (1 - 3r_{eD}^2)$$

$$W_{eD} = 0.5 (r_{eD}^2 - 1) (1 - \text{Exp}(C_2 t_D / J^*))$$

Due to the limited data available for analysis it was not possible to develop an analytical relationship between  $C_1$  and  $J^*$  or between  $C_2$  and  $t_D$ .

The main objective of this paper was to show the importance of varying the  $C_1$  and  $C_2$  constants based on the values of  $r_{eD}$  and  $t_D$  to conform with original solutions of Van Everdingen and Hurst. Consequently, it is recommended that future work should be devoted toward development of such analytical relations.

### Nomenclature

$c_f$	=	aquifer compressibility, $\text{psi}^{-1}$
$c_w$	=	water compressibility, $\text{psi}^{-1}$
$c_t$	=	$c_f + c_w$
$h$	=	aquifer thickness, ft
$k$	=	permeability, md
$r_o$	=	reservoir radius, ft
$r_e$	=	aquifer radius, ft
$r_{eD}$	=	dimensionless radius
$t$	=	time, year
$t_D$	=	dimensionless time
$\Phi$	=	porosity, fraction
$\mu_w$	=	water viscosity, cp

### REFERENCES

- [1] Dake, L.P. *Fundamentals Of Reservoir Engineering*.
- [2] Van Everdingen A.F. and Hurst W., 1949. The Application of Laplace Transformations to flow problems in reservoirs. *Trans. AIME*, **186**, 305-324.
- [3] Fetkovich M.J. A simplified approach to water influx calculations: finite aquifer systems *J. Pet. Tech.*, July, 814-828.
- [4] Marsal D., 1982. Topics of reservoir engineering. Course Notes Delft University of Technology.